## Tutorial 4

## February 16,2017

## 1. Discuss why there is no maximum principle for wave equation? Consider the following Cauchy problem:

$$\begin{cases} \partial_t^2 u - \partial_x^2 u = 0, & -\infty < x < +\infty, & t > 0\\ u(x, t = 0) = 0, & \partial_t u(x, t = 0) = \sin x, & -\infty < x < +\infty \end{cases}$$

And the unique solution is given by d'Alembert formula:

$$u(x,t) = \frac{1}{2}\cos(x+t) - \cos(x-t) = -\sin x \sin t, \quad -\infty < x < \infty, t > 0$$

Then u(x,t) attains its maximum 1 only at the interior points  $(\frac{\pi}{2}\pm 2n\pi, \frac{3\pi}{2}+2n\pi)$  or  $(\frac{3\pi}{2}\pm 2n\pi, \frac{\pi}{2}+2n\pi)$  for  $n = 0, 1, 2, \cdots$ . However, u(x,t) = 0 on the boundary  $\{(x,t) : t = 0\}$ . Therefore there is no maximum principle for the Cauchy problem for the 1-dimensitonal wave equation.

Remark: The key is to find an counterexample.

2. Use the Green's function of the heat equation to show that the backward heat equation is not wellposed.

Note that S(x,t) satisfies  $u_t = ku_{xx}$  for any t > 0, and  $S(0,t) \to \infty$  as  $t \to 0^+$ . Then u(x,t) = S(x,t+1) solves  $u_t = ku_{xx}$  for t > -1. Then  $S(0,t) \to \infty$  as  $t \to -1^+$ , which implies that there is no solution for the backward heat equation with initial data  $u(x,0) = S(x,1) = \frac{1}{\sqrt{4k\pi}}e^{-\frac{x}{4k}}$ , hence the backward heat equation is not well-posed.

3. Let  $\phi(x)$  be a continuous function such that  $|\phi(x)| \leq Ce^{ax^2}$ . Show that formula (8) on page 48 for the solution of the diffusion equation makes sense for  $0 < t < \frac{1}{4ak}$ , but not necessarily for larger t.

## Solution: Since

$$|e^{-(x-y)^2/4kt}\phi(y)| \le Ce^{-(x-y)^2/4kt+ay^2} = Ce^{(a-\frac{1}{4kt})y^2 + \frac{x}{2kt}y - \frac{x^2}{4kt}},$$
$$u(x,t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^{\infty} e^{-(x-y)^2/4kt}\phi(y) \, dy$$

makes sense for  $a - \frac{1}{4kt} < 0$ , i.e. 0 < t < 1/(4ak), but not necessarily for large t, for example,  $\phi(x) = e^{ax^2}$ .

4. Use energy method to show that the energy for diffution equation decays with a rate for large time. Multiplying  $\partial_t v = k \partial_x^2 v$  by v and then integrating w.r.t x give that

$$\frac{d}{dt}\int_{-\infty}^{\infty}\frac{1}{2}|v|^2dx = \int_{-\infty}^{\infty}k\partial_x^2vvdx$$

It follows from integration by parts that

$$\frac{d}{dt}\int_{-\infty}^{\infty}\frac{1}{2}|v|^2dx = k\partial_x vv\Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty}k(\partial_x v)^2dx = -\int_{-\infty}^{\infty}k(\partial_x v)^2dx \le 0$$

for any  $t \ge 0$ . Here assume that v vanishes when  $x \to \infty$ . Hence if the solution is not a constant,  $\frac{d}{dt} \int_{-\infty}^{\infty} \frac{1}{2} |v|^2 dx < 0$ , then the energy  $E = \int_{-\infty}^{\infty} \frac{1}{2} |v|^2 dx$  decays as  $t \to \infty$ .